



Mathematical
Institute

A Li-Yau and Aronson-Bénilan approach to study Keller-Segel

A. FERNÁNDEZ-JIMÉNEZ

*Mathematical Institute
University of Oxford*

Joint work with C. ELBAR & F. SANTAMBROGIO

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Aggregation of *Dictostelium discoideum* to create a more complex structure. Reference *Wikipedia*.

$$\begin{cases} \partial_t \rho = \Delta \rho^m - \operatorname{div}(\rho \nabla u) & \text{in } (0, \infty) \times \mathbb{R}^d, \\ -\Delta u = \rho & \text{in } (0, \infty) \times \mathbb{R}^d, \end{cases} \quad (\text{KS})$$

where $m = 2 - \frac{2}{d}$.

Goals:

- Dimension $d = 2$. A Li-Yau type estimate.
- Dimension $d \geq 3$. An Aronson-Bénilan type estimate.

Mass is preserved. $M := \int_{\mathbb{R}^2} \rho_0(x) dx = \int_{\mathbb{R}^2} \rho(x) dx.$

Second moment is **linear** [BDP06].

$$\frac{d}{dt} M_2(t) = \frac{d}{dt} \int_{\mathbb{R}^2} |x|^2 \rho(x) dx = 4M \left(1 - \frac{M}{8\pi}\right).$$

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➤ **Supercritical.** $M > 8\pi$, [HV96; RS14; CGMN22]

$$\rho(t) \rightarrow M\delta_0 + f \quad \text{as } t \nearrow T < \infty.$$

The so-called **chemotactic collapse**.

Using a HLS inequality on the energy we find a critical mass M_c .

- **Subcritical.** $M < M_c$, **diffusion** [Bed11; BCL09].
- **Critical.** $M = M_c$, globally well-posed and L^∞ norm globally bounded in time [BCL09].
- **Supercritical.** $M > M_c$, all radial solutions **blow up** in finite time [BK13].

➤ **Li-Yau** [LY86]. The heat equation $\partial_t \rho = \Delta \rho = \operatorname{div}(\rho \nabla(\log \rho))$,

$$\Delta(\log \rho(t, \cdot)) \geq -\frac{d}{2t}.$$

➤ **Aronson-Bénilan** [AB79]. PME $\partial_t \rho = \Delta \rho^m = \operatorname{div}(\rho \nabla(\frac{m}{m-1} \rho^{m-1}))$,

$$\Delta\left(\frac{m}{m-1} \rho^{m-1}(t, \cdot)\right) \geq -\frac{\alpha}{t}, \quad \alpha = \frac{d}{d(m-1) + 2}.$$

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☞ Aronson-Bénilan for **drift-diffusion** [KZ21; BPS23; DS24; DS25].

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☞ Li-Yau and Aronson-Bénilan for **non-local aggregation**-diffusion???

☞ L^∞ bounds on ρ in terms of $\delta := \inf_{x \in \mathbb{R}^d} \Delta v(x)$.

1. Differential inequality on the pressure.
2. L^∞ estimates.
3. Li-Yau and Aronson-Bénilan estimates.
4. Discussions and new implications.

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$$v(t, x) = \begin{cases} \log \rho(t, x) - u(t, x), & d = 2, \\ \frac{m}{m-1} \rho^{m-1}(t, x) - u(t, x), & d \geq 3, \quad \left(m = 2 - \frac{2}{d}\right). \end{cases}$$

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☞ **Differential inequality.**

$$\frac{d}{dt} \Delta v(t, x_0) \geq (\Delta v(t, x_0))^2 - \frac{(d-1)^2}{2d} \sup_{i,j} \|\partial_{ii} u - \partial_{jj} u\|_{L^\infty}^2$$

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☞ **Estimate** $\|\partial_{ii} u - \partial_{jj} u\|_{L^\infty}$.

$$\sup_{i,j} \|\partial_{ii} u - \partial_{jj} u\|_{L^\infty} \lesssim |\tilde{\delta}|^{d/(d+2)} \|\rho\|_{L^2}^{4/(d+2)},$$

where $\delta = \min(\Delta v)$ and $\tilde{\delta} = \max(\delta, \|\rho\|_{L^\infty})$

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Let us remember that

$$v(x) = \log \rho(x) - u(x) \quad \text{and} \quad \Delta v(x) = \Delta \log \rho(x) + \rho(x).$$

We want to study L^∞ estimates in terms of

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- **Small mass.** Assume $0 < M < \frac{8\pi}{e}$.
- **Subcritical mass.** Assume $0 < M < 8\pi$.
- **Critical mass.** Assume $M = 8\pi$.

Small mass

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☞ Arbitrary x_0 and $r > 0$

$$\log \rho(x_0) \leq \int_{B_r(x_0)} \log \rho + \frac{|\tilde{\delta}|}{8} r^2 \quad [\text{Subharmonic}]$$

$$\leq \log \left(\int_{B_r(x_0)} \rho \right) + \frac{|\tilde{\delta}|}{8} r^2 \leq \log \left(\frac{M}{\pi r^2} \right) + \frac{|\tilde{\delta}|}{8} r^2 \quad [\text{Jensen}]$$

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☞ **Optimal value** $r = \sqrt{8/|\tilde{\delta}|}$.

$$\log \rho(x_0) \leq \log(M|\tilde{\delta}|/8\pi) + 1 \quad \Rightarrow \quad \rho(x_0) \leq (Me/8\pi)|\tilde{\delta}|.$$

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☞ **L^∞ bound.** $\Delta v \geq \delta \Rightarrow \Delta \log \rho \geq \delta - \rho$

$$\|\rho\|_{L^\infty} \leq \frac{Me}{8\pi} |\delta - \|\rho\|_{L^\infty}| \quad \Rightarrow \quad \|\rho\|_{L^\infty} \leq \frac{Me}{8\pi - Me} |\delta|.$$

Assume $0 < M < 8\pi$ and $\delta \neq 0$. Let us argue by **contradiction**.

☞ **Sequence.** $\Delta \log \rho_n + \rho_n \geq \delta$ and $\|\rho_n\|_{L^\infty} \geq n\delta$.

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☞ **Scaling.** Consider

$$\eta_n(x) = \lambda_n^2 \rho_n(\lambda_n x) \quad \text{where} \quad \lambda_n^2 = \frac{1}{\|\rho_n\|_{L^\infty}}$$

and we notice

$$\int_{\mathbb{R}^2} \eta_n = M \quad \text{and} \quad \|\eta_n\|_{L^\infty} = 1.$$

Furthermore, $\Delta \log \eta_n + \eta_n = \Delta \log \rho_n(\lambda_n \cdot) + \lambda_n^2 \rho_n(\lambda_n \cdot) \geq \delta \lambda_n^2$.

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☞ **Compactness.** The sequence η_n converges to η such that

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Subsolution of the Liouville equation with mass $\leq M < 8\pi$.

Lemma (Subsolution of the Liouville equation)

Suppose $h: \mathbb{R}^2 \rightarrow [-\infty, \infty)$ is a function in $L^\infty(\mathbb{R}^2) \cap H_{loc}^1(\mathbb{R}^2)$ such that

$$\begin{cases} \Delta h + e^h \geq 0 & \text{in the set } \{ (x, y) : h(x, y) > t \} \text{ for all } t \in \mathbb{R}, \\ \int_{\mathbb{R}^2} e^h < +\infty. \end{cases}$$

Then $\int_{\mathbb{R}^2} e^h dx \geq 8\pi$.

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☞ **Contradiction.** Take $\eta = e^h$. Since $M < 8\pi$ contradiction.

☞ Due to the scaling $\|\rho\|_{L^\infty} \leq C|\delta|$.

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Example: dimension 3. Multiply by h , integrate and integrate by parts

$$\|\nabla h\|_{L^2}^2 \leq \frac{m-1}{m} \|h\|_{L^4}^4.$$

Gagliardo-Nirenberg. $\|h\|_{L^4}^4 \leq C_{GN} \|\nabla h\|_{L^2}^2 \|h\|_{L^3}^2 \leq C_{GN} \frac{m-1}{m} \|h\|_{L^3}^2 \|h\|_{L^4}^4.$

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- ➡ **Critical mass?** Analogous to dimension 2.

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Furthermore,
$$\|\rho(t, \cdot)\|_{L^\infty} \leq \frac{C}{t} \quad \text{for a.e. } (t, x) \in (0, T) \times \mathbb{R}^d$$

and we can construct solutions for $\rho_0 \in \mathcal{M}(\mathbb{R}^d)$ a nonnegative measure.

Recall $\delta' \geq \delta^2 - C \sup_{i,j} \|\partial_{ii} u - \partial_{jj} u\|_{L^\infty}^2$.

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Using **Log-HLS** and conservation of m_2 we get additional integrability.

$$\frac{d}{dt} \delta \geq \delta^2 (1 - o(1)) \quad \Rightarrow \quad \begin{cases} \delta \geq -C(T) \left(1 + \frac{1}{t}\right) \\ \|\rho\|_{L^\infty} \leq C(T) \left(1 + \frac{1}{t}\right) \end{cases}$$

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➤ **Subcritical mass** $d \geq 3$. $\mathcal{F}[\rho_0] < +\infty$. **HLS** to control $\int \rho^m$.

Recall $\delta' \geq \delta^2 - C \sup_{i,j} \|\partial_{ii} u - \partial_{jj} u\|_{L^\infty}^2$.

$\sup_{i,j} \|\partial_{ii} u - \partial_{jj} u\|_{L^\infty} \lesssim |\tilde{\delta}|^{d/(d+2)} \|\rho\|_{L^2}^{4/(d+2)}$ and $\tilde{\delta} = \max(\delta, \|\rho\|_{L^\infty})$.

➤ **Subcritical mass** $d = 2$. $\mathcal{F}[\rho_0], m_2[\rho_0] < +\infty$.

Using **Log-HLS** and conservation of m_2 we get additional integrability.

$$\frac{d}{dt} \delta \geq \delta^2 (1 - o(1)) \quad \Rightarrow \quad \begin{cases} \delta \geq -C(T) \left(1 + \frac{1}{t}\right) \\ \|\rho\|_{L^\infty} \leq C(T) \left(1 + \frac{1}{t}\right) \end{cases}$$

➤ **Subcritical mass** $d \geq 3$. $\mathcal{F}[\rho_0] < +\infty$. **HLS** to control $\int \rho^m$.

➤ **Critical mass**. The result can be extended.

➤ We construct solutions for Keller-Segel.

1. Differential inequality on the pressure.
2. L^∞ estimates.
3. Li-Yau and Aronson-Bénilan estimates.
4. Discussions and new implications.

➤ **Existence of solutions.** Uniform L^∞ bounds and global in time sol.

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- **Liouville and Lane-Emden.** Work in progress.

Thank you for you attention!!!

C. Elbar, A. F-J, and F. Santambrogio. A Li-Yau and Aronson-Bénilan approach for the Keller-Segel with critical exponent. *Preprint in preparation.*

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$$\begin{aligned}\partial_t \Delta v &= \Delta(\Delta v) + \nabla(\Delta v) \cdot [2\nabla v + \nabla u] + 2|D^2 v|^2 \\ &\quad + 2 \left(D^2 v - \frac{1}{2} \Delta v \text{Id} \right) : \left(D^2 u - \frac{1}{2} \Delta u \text{Id} \right)\end{aligned}$$

If $x_0 = x_0(t)$ is a minimum point of $\Delta v(t, \cdot)$

$$\begin{aligned}\partial_t \Delta v(t, x_0) &\geq 2 |D^2 v(t, x_0)|^2 \\ &\quad + 2 \left(D^2 v(t, x_0) - \frac{1}{2} \Delta v(t, x_0) \text{Id} \right) : \left(D^2 u(t, x_0) - \frac{1}{2} \Delta u(t, x_0) \text{Id} \right).\end{aligned}$$